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2D Computation and Measurement of Electric and Magnetic Fields of Overhead Electric Power Lines



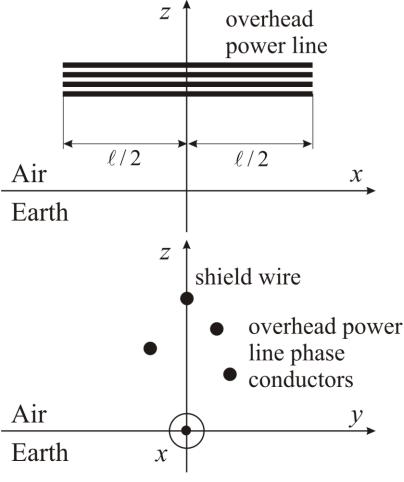
Introduction

- ➤ the problem of computing the power frequency electromagnetic field can be considered as quasistatic → electric and magnetic field computed separately
- computation of electric and magnetic field accomplished using a 2D approximation of a power line
- short power line is considered, conductor charge density is approximated by a constant
- results will be compared to measurements underneath a 400 kV overhead power line



Mathematical model of overhead electric power line

- > 2D numerical algorithm → short conductors
 → conductor sag neglected
 > conductors positioned along *x*-axis
 > conductors treated as line sources parallel to earth surface
 - \succ thin-wire approximation
 - ➢ homogeneous earth
 - Field distribution will be computed along y-z plane in the middle of the section





Scalar electric and vector magnetic potentials

- solutions of Helmholtz differential equations reduce to solutions of Poisson
 differential equations (quasistatic approximation)
- \succ scalar electric potential for *n* conductors:

$$\overline{\varphi} = \frac{1}{4 \cdot \pi \cdot \varepsilon_0} \cdot \sum_{i=1}^n \left[\int_{\Gamma_{Ai}} \frac{\overline{\lambda}_i}{R_{Ai}} \cdot d\ell_{Ai} + \overline{F} \cdot \int_{\Gamma_{Bi}} \frac{\overline{\lambda}_i}{R_{Bi}} \cdot d\ell_{Bi} \right]$$

 \blacktriangleright vector magnetic potential for *n* conductors:

$$\underline{\vec{A}} = \vec{i} \cdot \frac{\mu_0}{4 \cdot \pi} \cdot \sum_{i=1}^n \bar{I}_i \cdot \int_{\Gamma_{Ai}} \frac{1}{R_{Ai}} \cdot d\ell_{Ai}$$



- > numerical algorithm is based on **Biot-Savart** law
- \triangleright conductor current \bar{I}_i flows along the axis and generates the magnetic field
- \triangleright components of magnetic flux density in *y*-*z* plane for x = 0:

$$\overline{B}_{x} = 0$$

$$\overline{B}_{y} = \mu_{0} \cdot \sum_{i=1}^{n} \frac{(z_{i} - z) \cdot \ell}{4 \cdot \pi \cdot d_{i}^{2} \cdot \sqrt{d_{i}^{2} + \frac{\ell^{2}}{4}}} \cdot \overline{I}_{i}$$

$$\overline{B}_{z} = \mu_{0} \cdot \sum_{i=1}^{n} \frac{(y - y_{i}) \cdot \ell}{4 \cdot \pi \cdot d_{i}^{2} \cdot \sqrt{d_{i}^{2} + \frac{\ell^{2}}{4}}} \cdot \overline{I}_{i}$$

$$\overline{B}_{z} = \mu_{0} \cdot \sum_{i=1}^{n} \frac{(y - y_{i}) \cdot \ell}{4 \cdot \pi \cdot d_{i}^{2} \cdot \sqrt{d_{i}^{2} + \frac{\ell^{2}}{4}}} \cdot \overline{I}_{i}$$



> conductor charge density $\overline{\lambda}_i$ approximated by a **constant** $\rightarrow \overline{\lambda}_i = \frac{Q_i}{\ell}$ > electric field intensity is computed from

$$\underline{\vec{E}} = -\nabla\overline{\varphi} - j \cdot \omega \cdot \underline{\vec{A}} = \left\{ \overline{E}_x, \overline{E}_y, \overline{E}_z \right\}$$

➤ the components of the electric field intensity are computed from:

$$\begin{split} \overline{E}_{x} &= \overline{E}_{x} \big|_{x=0} = -j \cdot \omega \cdot \overline{A}_{x} = -j \cdot \omega \cdot \overline{A} \big|_{x=0} & \text{rms value:} \\ \overline{E}_{y} &= \overline{E}_{y} \big|_{x=0} = -\frac{\partial \overline{\phi}(x, y, z)}{\partial y} \big|_{x=0} = -\frac{\partial \overline{\phi}(0, y, z)}{\partial y} & E = \sqrt{E_{x} + E_{y} + E_{z}} \\ \overline{E}_{z} &= \overline{E}_{z} \big|_{x=0} = -\frac{\partial \overline{\phi}(x, y, z)}{\partial z} \big|_{x=0} = -\frac{\partial \overline{\phi}(0, y, z)}{\partial z} \end{split}$$



 \blacktriangleright vector magnetic potential for x = 0 needed in:

$$\overline{E}_{x} = -j \cdot \omega \cdot \overline{A} \Big|_{x=0}$$

is computed from:
$$\overline{A} \Big|_{x=0} = \frac{\mu_{0}}{2 \cdot \pi} \cdot \sum_{i=1}^{n} \overline{I}_{i} \cdot \ell n \frac{\sqrt{d_{i}^{2} + \frac{\ell^{2}}{4}} + \frac{\ell}{2}}{d_{i}} - \frac{\frac{\Gamma_{Ai}}{2}}{-\frac{\ell}{2}} \frac{\Gamma_{Ai}}{x \, dx} \frac{\ell}{\frac{\ell}{2}} x$$

$$d_{i} = \sqrt{(y-y_{i})^{2} + (z-z_{i})^{2}} \qquad \qquad \int_{\Gamma_{Ai}} \frac{1}{R_{Ai}} \cdot d\ell_{Ai}$$



 \blacktriangleright scalar electric potential for x = 0 needed in equations

$$\overline{E}_{y} = -\frac{\partial \varphi(0, y, z)}{\partial y}$$
 $\overline{E}_{z} = -\frac{\partial \varphi(0, y, z)}{\partial z}$

can be obtained from:

$$\overline{\varphi}(0, y, z) = \frac{1}{2 \cdot \pi \cdot \varepsilon_0 \cdot \ell} \cdot \sum_{i=1}^n \left[\sinh^{-1} \left(\frac{\ell}{2 \cdot d_i} \right) + \overline{F} \cdot \sinh^{-1} \left(\frac{\ell}{2 \cdot D_i} \right) \right] \cdot \overline{Q}_i$$

 \succ phasors of the *i*th conductor charges are unknown



phasors of the *i*th conductor charges are computed using the point collocation method:

$$\sum_{i=1}^{n} \left[\overline{Z}(d_{ji}) + \overline{F} \cdot \overline{Z}(D_{ji}) \right] \cdot \overline{Q}_{i} = \overline{\Phi}_{j} \quad ; \quad j = 1, 2, ..., n$$
$$\overline{Z}(v) = \frac{1}{2 \cdot \pi \cdot \varepsilon_{0} \cdot \ell} \cdot \ell n \frac{\sqrt{\left(\frac{\ell}{2}\right)^{2} + v^{2}} + \frac{\ell}{2}}{v}$$

 \succ *v* is replaced by:

$$d_{ii} = r_{0i} \qquad d_{ij} = \sqrt{(y_i - y_j)^2 + (z_i - z_j)^2} \quad ; \quad i \neq j$$
$$D_{ii} = 2 \cdot z_i \qquad D_{ij} = \sqrt{(y_i - y_j)^2 + (z_i + z_j)^2} \quad ; \quad i \neq j$$



Input data of the overhead power line

- 6 phase conductors (AlFe 490/65) and 2 shield wires (Alumweld 19/9)
- ➢ 376 m long power line section

Z							
Shield wire 1	Shield wire 2						
	• L3 •						

У

	i	y (m)	z_{s} (m)	$\overline{\Phi}_i$ (kV)	\bar{I}_i (A)
L1	1	-10.7	12.58	234.4∠0°	228∠-4.2°
	2	-10.4	12.58	234.4∠0°	228∠-4.2°
L2	3	-0.15	12.58	234.4∠240°	228∠235.8°
	4	0.15	12.58	234.4∠240°	228∠235.8°
L3	5	10.4	12.58	234.4∠120°	228∠115.8°
	6	10.7	12.58	234.4∠120°	228∠115.8°
SW1	7	-7.44	19.09	0∠0°	0∠0°
SW2	8	7.44	19.09	0∠0°	0∠0°



Input data of the overhead power line

The sag of the conductors was taken into account by:

$$z_s = z - \frac{2}{3} \cdot s$$

- measurements taken on a clear winter day (10 °C)
- terrain mostly clear except sporadic
 bushes
- measurements along y directed observation profile at z = 1 m

Other input data:

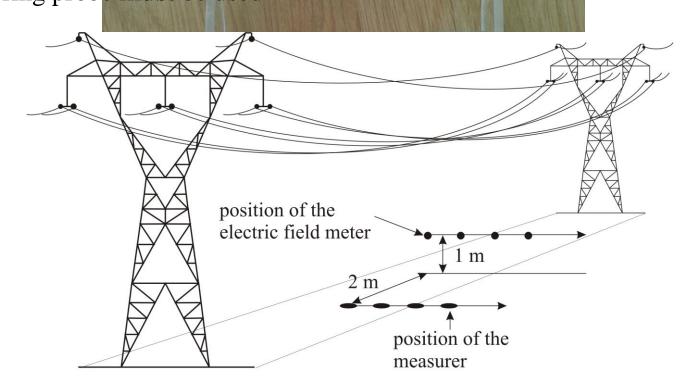
 $\varepsilon_r = 10$ $\sigma = 0.1 S / m$ f = 50 Hz





Measuring equipment and procedure

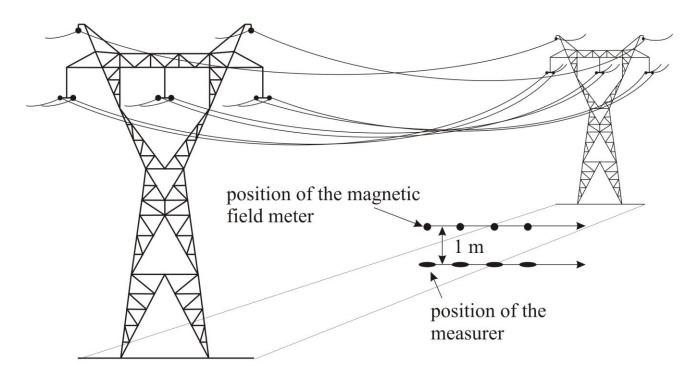
- electric field meter Monroe Electronics Model 238A-1
 AC Fieldmeter
- ➤ measurer has influence on electric field distribution → measuring probe must be used





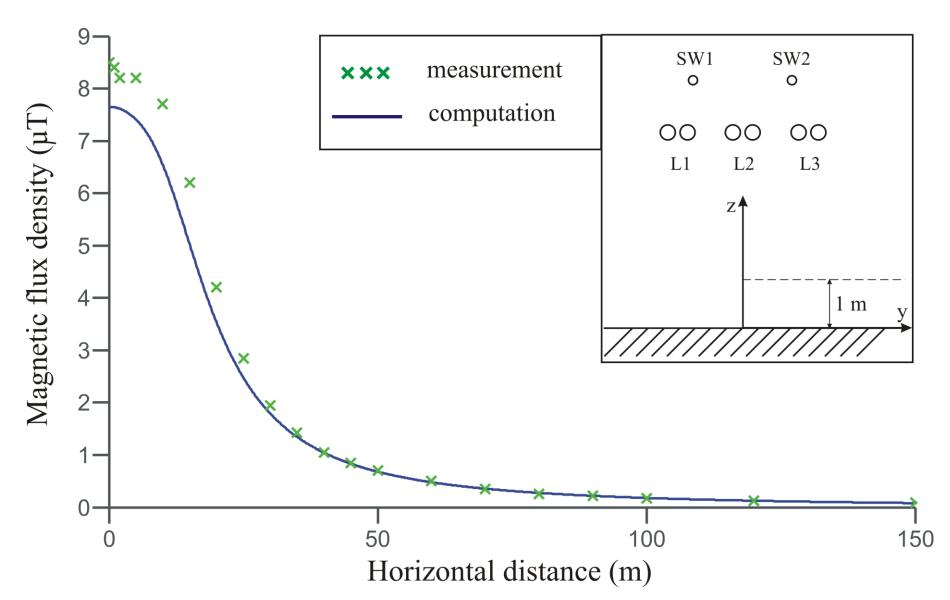
Measuring equipment and procedure

- magnetic field meter Sypris Triaxial ELF Magnetic Field Meter Model 4090
- measurer does not have influence on the magnetic field distribution



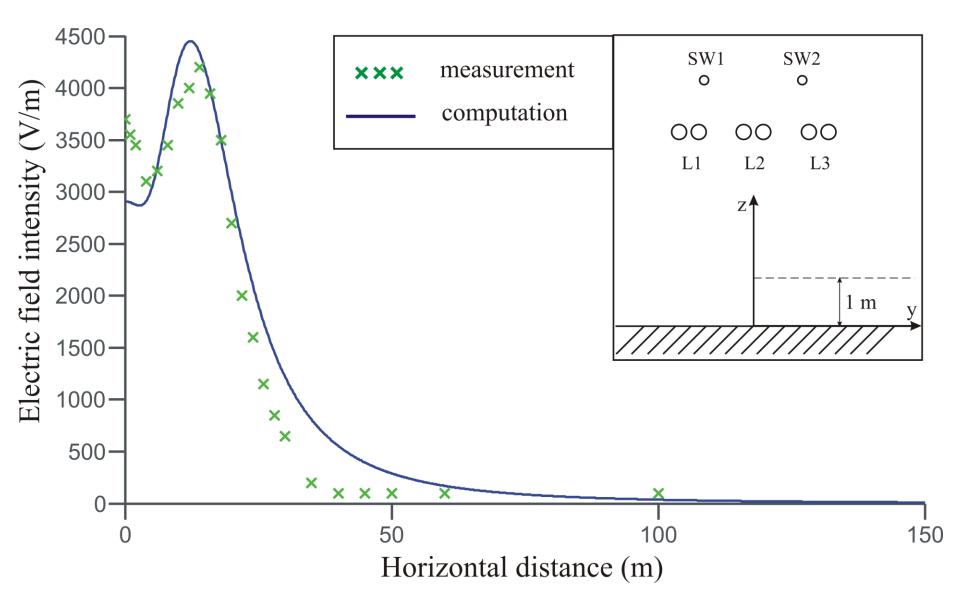


Comparison of results – magnetic field





Comparison of results – electric field





Summary

- 2D algorithms for computation of electric and magnetic fields of overhead power lines
- > good agreement was found between the computed results and measurements
- ➢ better results can be achieved by taking the sag into account (3D algorithm)
- more advanced models are needed for more complicated structures such as electric power substations



Thank you!